

On the Feynman graph expansion of 1-particle irreducible n-point functions in quantum field theory

Ângela Mestre

Institut de Minéralogie et de Physique des Milieux Condensés, Paris

March 9, 2010

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- Field operator algebra

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- Field operator algebra
- Algebraic representation of graphs
- Coalgebra structures
- Linear maps
- Graph generation and applications to QFT

Field operator algebra [Brouder & Oeckl 2003, Brouder 2009]

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- antipode:

$$S(\phi(x_1) \dots \phi(x_n)) := (-1)^n \phi(x_1) \dots \phi(x_n).$$

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$$\int dx dy dz dx' dy' dz'$$

$$G_F(x_1, x) \Sigma^{(3)}(x, y, z) G_F(y, y') G_F(z, z') \Sigma'^{(3)}(x', y', z') G_F(x', x_2)$$

$$= \int dy dz dy' dz' \mathcal{V}^{(3)}(x_1, y, z) G_F^{-1}(y, y') G_F^{-1}(z, z') \mathcal{V}^{(3)}(x_2, y', z').$$

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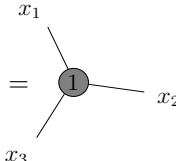
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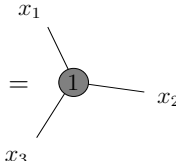
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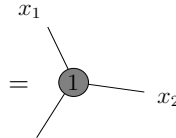
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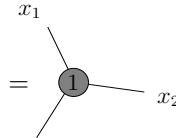
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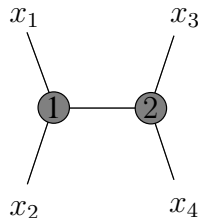
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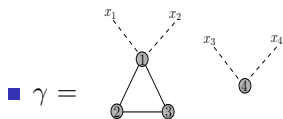
$$S_{1,\dots,v}^{\gamma} := \text{int}_{1,\dots,v}^{\gamma} \cdot \text{ext}_{1,\dots,v}^{\gamma}.$$

Examples

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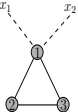

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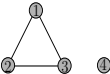
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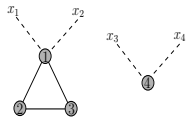



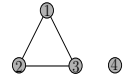
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$$S_{\sigma(1), \dots, \sigma(v)}^{\gamma} := \text{int}_{\sigma(1), \dots, \sigma(v)}^{\gamma} \cdot \text{ext}_{\sigma(1), \dots, \sigma(v)}^{\gamma},$$

where

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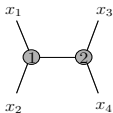
Examples

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- $S_{1,2}^\gamma \in S(V)^{\otimes 2} =$

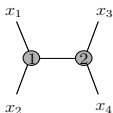
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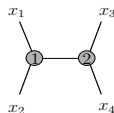
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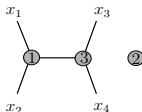
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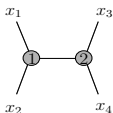


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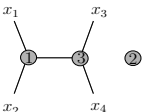


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$$(\sigma : \{1, 2\} \rightarrow \{1, 3\}; 1 \mapsto 1, 2 \mapsto 3)$$

A non-associative non-commutative graph multiplication

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$$\begin{aligned} (u_1 \otimes \dots \otimes u_v) \bullet_{i,j} (u'_1 \otimes \dots \otimes u'_{v'}) &:= \\ (\tau_{v-1} \circ \dots \circ \tau_i)(u_1 \otimes \dots \otimes u_v) \otimes (\tau_2 \circ \dots \circ \tau_j)(u'_1 \otimes \dots \otimes u'_{v'}) & \\ = u_1 \otimes \dots \otimes \hat{u}_i \otimes \dots \otimes u_v \otimes u_i \otimes u'_j \otimes u'_1 \otimes \dots \otimes \hat{u}'_j \otimes \dots \otimes u'_{v'} . & \end{aligned}$$



$$S_{1,\dots,v}^{\gamma} \bullet_{i,j} S_{1,\dots,v'}^{\gamma'} := S_{\sigma(1),\dots,\sigma(v)}^{\gamma} \cdot S_{\sigma'(1),\dots,\sigma'(v')}^{\gamma'}$$

$$(S_{\sigma(1),\dots,\sigma(v)}^{\gamma} \cdot S_{\sigma'(1),\dots,\sigma'(v')}^{\gamma'}) = \text{disconnected graph}$$

Gluing two graphs at a vertex:

$$\diamond_{i,j} := \tau_{v+v'-2} \circ \dots \circ \tau_v \circ \cdot_v \circ \bullet_{i,j} : S(V)^{\otimes v} \times S(V)^{\otimes v'} \rightarrow S(V)^{\otimes (v+v'-1)}$$

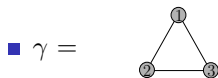
$(v + v' - 1 = \text{cut vertex})$

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
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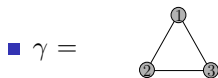


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
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Examples




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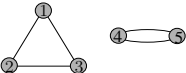
■ $\gamma' =$ 

■ $\gamma \bullet_{3,2} \gamma' =$


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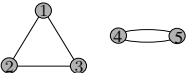
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
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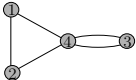
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A coalgebra structure on 1PI graphs

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$\mathfrak{B}_{l,v}$ = set of all 1PI Feynman graphs on l loops, v vertices and no external edges nor self-loops.

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$$\Delta(\bar{s}) := \bar{s} \otimes \bar{s};$$

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$$\begin{aligned} \Delta(\mathbf{1}) &:= \mathbf{1} \otimes \mathbf{1}; \\ \Delta(B_{1,\dots,v}^{\gamma}) &:= \frac{1}{v} \sum_{i=1}^v \Delta_i(B_{1,\dots,v}^{\gamma}), \end{aligned}$$

where

$$\begin{aligned} \Delta_i(B_{1,\dots,v}^{\gamma}) &:= B_{\sigma_i(1),\dots,\sigma_i(v)}^{\gamma} + B_{\sigma_{i+1}(1),\dots,\sigma_{i+1}(v)}^{\gamma} \\ &= B_{1,\dots,\widehat{i+1},i+2,\dots,v+1}^{\gamma} + B_{1,\dots,\hat{i},i+1,\dots,v+1}^{\gamma}. \end{aligned}$$

■ counit $\epsilon : \mathcal{B} \rightarrow \mathbb{C}$:

$$\epsilon(\mathbf{1}) := 1;$$

$$\epsilon(\mathcal{B}_{l,v}) := 0 \quad \text{if } v > 1.$$

Generalizing to connected graphs

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Extend $\Delta := \frac{1}{v} \sum_{i=1}^v \Delta_i$ to \mathcal{B}^* by requiring the maps Δ_i to satisfy:

$$\Delta_i \left(\prod_{a=1}^k B_{\sigma_a(1), \dots, \sigma_a(v_a)}^{\gamma_a} \right) := \prod_{a=1}^k \Delta_i \left(B_{\sigma_a(1), \dots, \sigma_a(v_a)}^{\gamma_a} \right).$$

counit $\epsilon : \mathcal{B}^* \rightarrow \mathbb{C}$:

$$\epsilon(1) := 1;$$

$$\epsilon\left(\prod_{a=1}^k B_{\sigma_a(1), \dots, \sigma_a(v_a)}^{\gamma_a}\right) := 0 \quad \text{if } k > 0.$$

Maps $Q_{i \geq 1}^{(\rho)}$ and $\hat{Q}_{i \geq 1}^{(\rho)}$

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Truncated coproduct: $\Delta_{\geq 1} : V^n \rightarrow \bigoplus_{i=1}^{n-1} V^i \otimes V^{n-i}$ [M. & Oeckl 2006]:

$$\begin{aligned} \Delta_{\geq 1}(1) &= 0, & \Delta_{\geq 1}(\phi(x)) &= 0; \\ \Delta_{\geq 1}(\phi(x)\phi(y)) &= \phi(x) \otimes \phi(y) + \phi(y) \otimes \phi(x). \end{aligned}$$

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For all $1 \leq i \leq v$, define:

$$Q_{i \geq 1}^{(\rho)} := \frac{1}{2(\rho-1)!} R_{i, i+1}^{\rho} \cdot \Delta_{i \geq 1} : S(V)^{\otimes v} \rightarrow S(V)^{\otimes v+1}.$$

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$(Q_{i \geq 1}^{(1)})$ [Glover *et al* 1979], [Livernet 2006])

Examples:

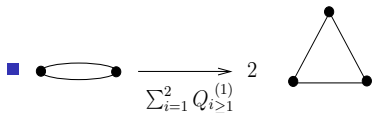
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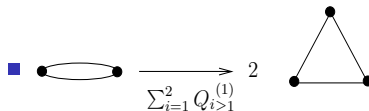
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On \mathcal{B}^* :

$$\hat{Q}_{i \geq 1}^{(\rho)} \left(\prod_{a=1}^k B_{\sigma_a(1), \dots, \sigma_a(v_a)}^{\gamma_a} \right) := \frac{1}{2(\rho - 1)!} R_{i, i+1}{}^\rho \prod_{a=1}^k \Delta_{i \geq 1} (B_{\sigma_a(1), \dots, \sigma_a(v_a)}^{\gamma_a}).$$

Maps $B_{\pi_i(1), \dots, \pi_i(v)}^\gamma \cdot \Delta_i^{v-1}$

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Define:

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Examples

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$$\begin{aligned}
 R_{3,4} \cdot \Delta_3(B_{1,2,3}^{C_3} \cdot B_{3,4,5}^{C_3}) &= R_{3,4} \cdot \Delta_3(B_{1,2,3}^{C_3}) \cdot \Delta_3(B_{3,4,5}^{C_3}) \\
 &= R_{3,4} \cdot (B_{1,2,3}^{C_3} + B_{1,2,4}^{C_3}) \cdot (B_{3,5,6}^{C_3} + B_{4,5,6}^{C_3}) \\
 &= R_{3,4} \cdot B_{1,2,3}^{C_3} \cdot B_{3,5,6}^{C_3} + R_{3,4} \cdot B_{1,2,3}^{C_3} \cdot B_{4,5,6}^{C_3} + \\
 &\quad R_{3,4} \cdot B_{1,2,4}^{C_3} \cdot B_{3,5,6}^{C_3} + R_{3,4} \cdot B_{1,2,4}^{C_3} \cdot B_{4,5,6}^{C_3};
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 &= R_{3,4} \cdot B_{1,2,3}^{C_3} \cdot B_{3,5,6}^{C_3} + R_{3,4} \cdot B_{1,2,3}^{C_3} \cdot B_{4,5,6}^{C_3} + \\
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 \end{aligned}$$

$R_{3,4} \cdot \Delta_3$ (triangle with internal vertex) = 2 (triangle with internal vertex and external line) + 2 (two triangles connected by a line)

Generating 1PI Feynman graphs

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Pick out the terms that generate 1VI graphs according to a formula given in [M. 2009]:

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Theorem (see M. 2010)

For all integers $l \geq 0$ and $v > 1$, define $\mathfrak{Y}_{1,\dots,v}^{l,v} \in S(V)^{\otimes v}$ by the following recursion relation:

- $\mathfrak{Y}_{1,2}^{l,2} := \frac{1}{2^{(l+1)!}} R_{1,2}^{l+1};$
- $\mathfrak{Y}_{1,\dots,v}^{0,v} := 0, v > 2;$
-

$$\mathfrak{Y}_{1,\dots,v}^{l,v} := \frac{1}{l+v-1} \left(\sum_{\rho=1}^{l+1} \sum_{i=1}^{v-1} Q_{i \geq 1}^{(\rho)} (\mathfrak{Y}_{1,\dots,v-1}^{l+1-\rho, v-1}) + \sum_{j=2}^{v-2} \sum_{\rho=1}^{l-j+1} \hat{Q}_{v-1 \geq 1}^{(\rho)} (\mathfrak{Y}_{1,\dots,v-1}^{l+1-\rho, v-1, j}) \right),$$

Theorem (cont.)

where for all integers $j > 1$, $v \geq j + 1$ and $l \geq j$, $\mathfrak{B}_{1,\dots,v}^{l,v,j}$ is given by the following recursion relation:

$$\mathfrak{B}_{1,\dots,v}^{l,v,2} := \frac{1}{l+v-1} \sum_{l'=1}^{l-1} \sum_{v'=2}^{v-1} \sum_{i=1}^{v'} \sum_{j=1}^{v-v'+1} \left((l' + v' - 1) \mathfrak{B}_{1,\dots,v'}^{l',v'} \diamond_{i,j} \mathfrak{B}_{1,\dots,v-v'+1}^{l-l',v-v'+1} \right);$$

$$\mathfrak{B}_{1,\dots,v}^{l,v,j} := \frac{1}{l+v-1} \sum_{l'=1}^{l-1} \sum_{v'=2}^{v-1} \sum_{i=1}^{v'-1} \left((l' + v' - 1) \mathfrak{B}_{1,\dots,v'}^{l',v'} \diamond_{i,v-v'+1} \mathfrak{B}_{1,\dots,v-v'+1}^{l-l',v-v'+1,j-1} \right).$$

Then, for fixed values of v and l , $\mathfrak{B}_{1,\dots,v}^{l,v}$ is the weighted sum over all 1VI Feynman graphs with l loops, v vertices and no external edges nor self-loops, each with weight given by the inverse of its symmetry factor.

Generating 1PI Feynman graphs

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Theorem (see M. 2010)

For all integers $l > 0$ and $v > 1$, define $\mathfrak{J}_{1,\dots,v}^{v,l} \in S(V)^{\otimes v}$ by the following recursion relation:

- $\mathfrak{J}_{1,2}^{l,2} := \mathfrak{W}_{1,2}^{l,2};$

-

$$\mathfrak{J}_{1,\dots,v}^{l,v} := \mathfrak{W}_{1,\dots,v}^{l,v} + \frac{1}{l+v-1} \cdot \sum_{l'=1}^{l-1} \sum_{v'=2}^{v-1} \sum_{i=1}^{v-v'+1} \left((l'+v'-1) \mathfrak{W}_{\pi_i(1),\dots,\pi_i(v)}^{l',v'} \cdot \Delta_i^{v-1}(\mathfrak{J}_{1,\dots,v-v'+1}^{v-l'+1,l-l'}) \right), v > 2.$$

Then, for fixed values of v and l , $\mathfrak{J}_{1,\dots,v}^{l,v}$ is the weighted sum over all 1PI Feynman graphs with l loops, v vertices and no external edges nor self-loops, each with weight given by the inverse of its symmetry factor.

Example of calculation

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$$\begin{aligned}
 & \frac{1}{2 \cdot 10} \left(\text{diagram} \cdot \Delta_{\text{cutvertex}} \left(\frac{1}{8} \right) + \text{diagram} \cdot \sum_{i \neq * } \Delta_i \left(\frac{1}{4} \right) \left(\text{diagram}^* \right) \right) = \\
 & = \frac{1}{2 \cdot 10} \left(\left(\frac{1}{2} + \frac{3}{4} \right) \text{diagram} + \left(1 + \frac{3}{2} \right) \text{diagram} + \left(\frac{1}{2} + \frac{3}{4} \right) \text{diagram} \right) + \dots \\
 & = \frac{1}{24} \text{diagram} + \frac{1}{23} \text{diagram} + \frac{1}{24} \text{diagram} + \dots
 \end{aligned}$$

External edges and self-loops

External edges and self-loops

- $T_i := \frac{1}{2}R_{i,i} : S(V)^{\otimes v} \rightarrow S(V)^{\otimes v}$ with $1 \leq i \leq v$ [M. & Oeckl 2006];

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Proposition

Fix an integer $n \geq 0$ as well as operator labels x_1, \dots, x_n . For all integers $l \geq 1$, $l' \geq 0$ and $v \geq 1$, define $\Gamma^{l+l',v} : S(V) \rightarrow S(V)^{\otimes v}$ as follows:

$$\begin{aligned} \Gamma^{l',1} &:= \frac{1}{l'!} T_1^{l'}; \\ \Gamma^{l+l',v} &:= \frac{1}{l'!} \mathfrak{J}_{1,\dots,v}^{l',v} \cdot \delta^{v-1}(T_1^{l'}) \cdot \Delta^{v-1}, v \geq 2, \end{aligned}$$

Then, $\Gamma^{l+l',v}(\phi(x_1) \cdots \phi(x_n))$ is the weighted sum over all 1PI Feynman graphs with l loops, l' self-loops, v vertices and n external edges whose end points are labeled x_1, \dots, x_n , each with weight given by the inverse of its symmetry factor.

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Contributions to the ensemble τ of 1PI n -point functions:

$$\tau^{l+l'} = \sum_{v=1}^{\infty} \tau^{l+l',v}, \quad \tau = \sum_{l+l'=0}^{\infty} \tau^{l+l'}.$$

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Vertex order contributions:

Corollary

For $v \geq 1$:

$$\tau^{l+l',v} = \nu_{1PI}^{\otimes v} \circ \Gamma^{l+l',v}.$$

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